

Action Potential Propagation in Branched Cable Network

Adv Topics Math Physiology: Modeling Neuronal
Dynamics

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Overview

- **Equations: Cable Equations (PDE) and Hodgkin Huxley Kinetics**
- **Construction of tree data structure**
- **Implement algorithm that solves tridiagonal system on the tree**
- **Experiment to see whether propagation fails or succeeds at a junction (all or none)**
- **Can we get differential propagation with this model ?**

Cable Equations with Hodgkin Huxley Kinetics:

Interior Equations:

$$c \frac{\partial v}{\partial t} + g(v-E) = \frac{r}{z_0} \frac{\partial^2 v}{\partial x^2}$$

$$g = g_{Na} + g_K + g_L$$

$$E = \frac{g_{Na} E_{Na} + g_K E_L + g_L E_L}{g_{Na} + g_K + g_L}$$

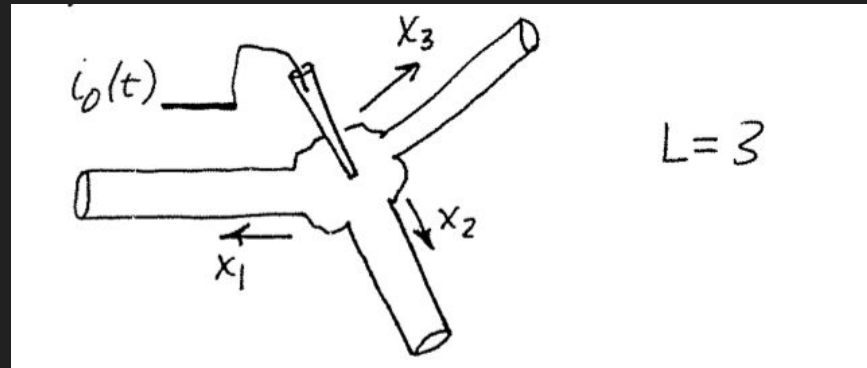
$$g_{Na} = \bar{g}_{Na} m^3 h$$

$$g_K = \bar{g}_K n^4$$

$$g_L = \bar{g}_L$$

$$\frac{\partial s}{\partial t} = \alpha_s(v)(1-s) - \beta_s(v)s$$

where $s = m, h, \text{ or } n$



Boundary Condition: Voltage is continuous and current adds up

$$v_0(t) = v_1(0,t) = v_2(0,t) = \dots = v_L(0,t)$$

$$i_0(t) = A_0 \left(c \frac{\partial v_0}{\partial t} + g_0(v_0 - E_0) \right) + \sum_{l=1}^L \left(-\frac{\pi r_l^2}{\rho} \right) \frac{\partial v_l}{\partial x_l}(0,t)$$

Numerical Method: Crank-Nicolson Method

$$\frac{S_j^{k+1/2} - S_j^{k-1/2}}{\Delta t} = \alpha_s(v_j^k) \left(1 - \frac{S_j^{k+1/2} + S_j^{k-1/2}}{2} \right) - \beta_s(v_j^k) \left(\frac{S_j^{k+1/2} + S_j^{k-1/2}}{2} \right)$$

$$C \frac{v_j^{k+1} - v_j^k}{\Delta t} + \tilde{g}_j^{k+1/2} \left(\frac{v_j^{k+1} + v_j^k}{2} - E_j^{k+1/2} \right) = \frac{r}{2\rho} \left(D^+ D^- \frac{v^{k+1} + v^k}{2} \right)_j$$

$$i_0^{k+1/2} = \tilde{A}_0 \left(C \frac{v_0^{k+1} - v_0^k}{\Delta t} + \tilde{g}_0^{k+1/2} \left(\frac{v_0^{k+1} + v_0^k}{2} - E_0^{k+1/2} \right) \right) + \sum_{l=1}^L \left(-\frac{\pi r_l^2}{\rho} \right) \frac{\left(\frac{v_{l1}^{k+1} + v_{l1}^k}{2} \right) - \left(\frac{v_0^{k+1} + v_0^k}{2} \right)}{\Delta x_l}$$

$$\tilde{A}_0 = A_0 + \sum_{l=1}^L \frac{1}{2} (2\pi r_l \Delta x_l)$$

$$\tilde{g}_0 = \frac{A_0 g_0 + \sum_{l=1}^L \frac{1}{2} (2\pi r_l \Delta x_l) g_{l0}}{\tilde{A}_0}$$

$$\tilde{E}_0 = \frac{A_0 g_0 E_0 + \sum_{l=1}^L \frac{1}{2} (2\pi r_l \Delta x_l) g_{l0} E_{l0}}{\tilde{A}_0 \tilde{g}_0}$$

Construction of Tree data structure:

- Arbitrary connected graph of nodes with no loops.
- Number the nodes $1, \dots, N$ and arbitrarily select one to be the root R .
- Since G is a connected tree, it is possible to find a unique chain between 1 and any node i .
- $R(i)$ is the root-ward neighbor of node i . Note $R(1) = 0$.
- $fc(i) = j$ is the leafward neighbor of node i with the smallest index (first child)
- If i has multiple leaf-ward neighbors (e.g. j, p, q), then these are denoted as “next sibling” of the “child” before it : $ns(j) = p$, $ns(p) = q$, $ns(q) = 0$ if its the last sibling.

$$-a_i \cdot V_{R(i)} + b_i V_i - \sum_{j \in \mathcal{L}(i)} c_j V_j = W_j$$

Simple tree example:

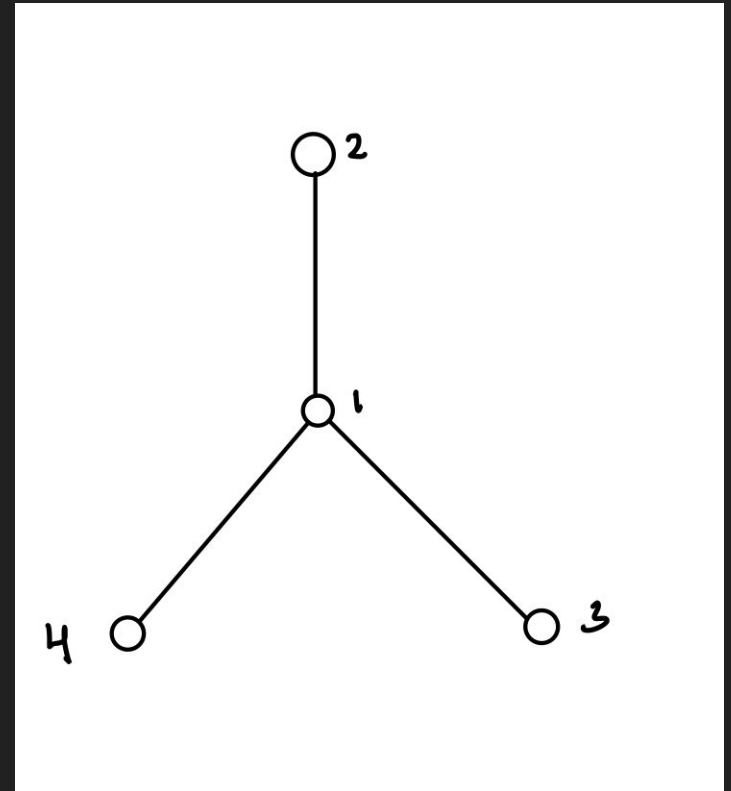
$$fc(1) = 2 \quad ns(1) = 0 \quad R(1) = 0$$

$$fc(2) = 0 \quad ns(2) = 1 \quad R(2) = 1$$

$$fc(3) = 0 \quad ns(3) = 4 \quad R(3) = 1$$

$$fc(4) = 0 \quad ns(4) = 0 \quad R(4) = 1$$

This is a high-level description of the basic tree
tree will run simulations on, however, more nodes
will be added along the branches.



Tridiagonal Solver Along the Tree:

```
function v = vnew(a, b, c, W, R, fc, ns, J)
v=W;

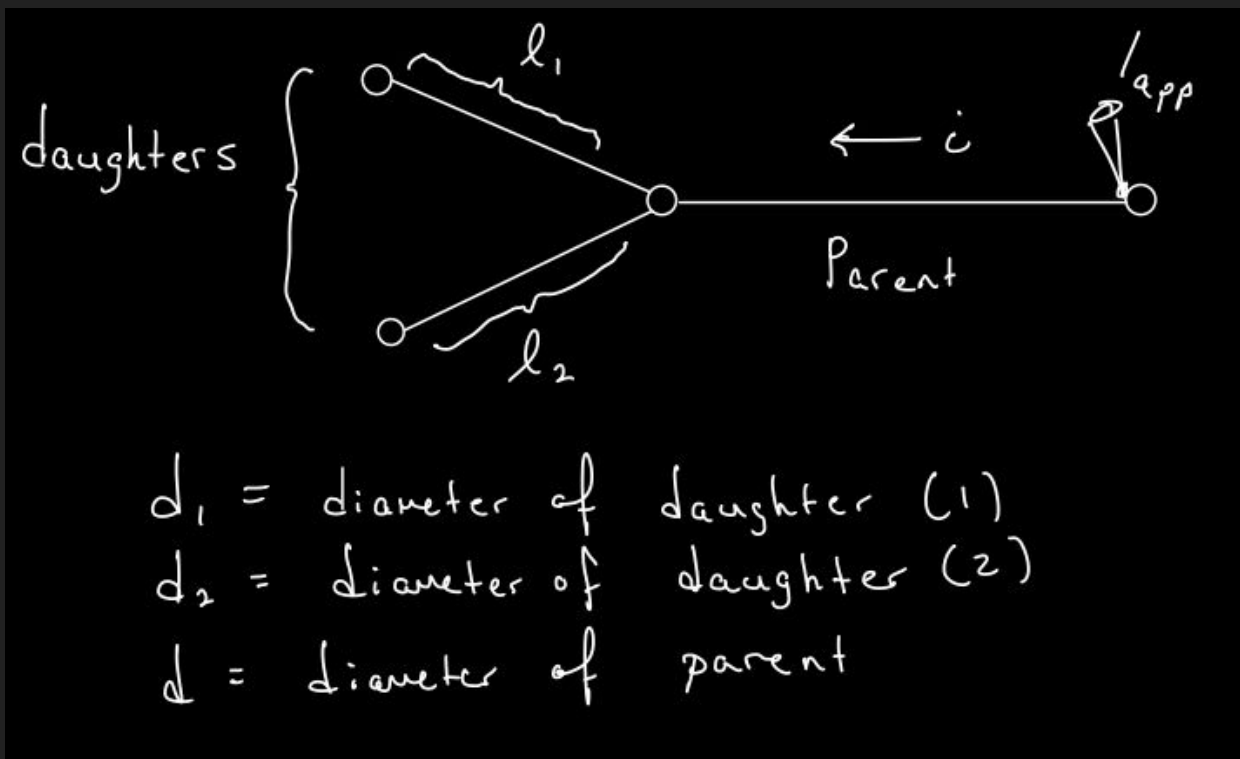
for i=J:(-1):1
    j=fc(i);
    while(j~=0)
        b(i)=b(i)-c(j)*a(j);
        v(i)=v(i)+c(j)*v(j);
        j=ns(j);
    end
    a(i)=a(i)/b(i);
    v(i)=v(i)/b(i);
end
for i=2:(+1):J
    v(i) = v(i) + a(i)*v(R(i));
end
```

Colliding AP's

Colliding action potentials cancel each other out because the refractory period of either spike prevents the continuation of an impulse in either direction:

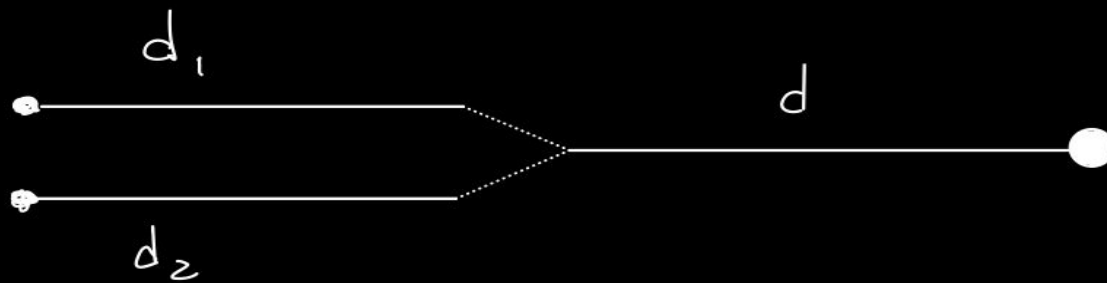


Consider a parent leading into two daughters:



Rall's Equivalent Cylinder (1962)

Let $l_1 = l_2$

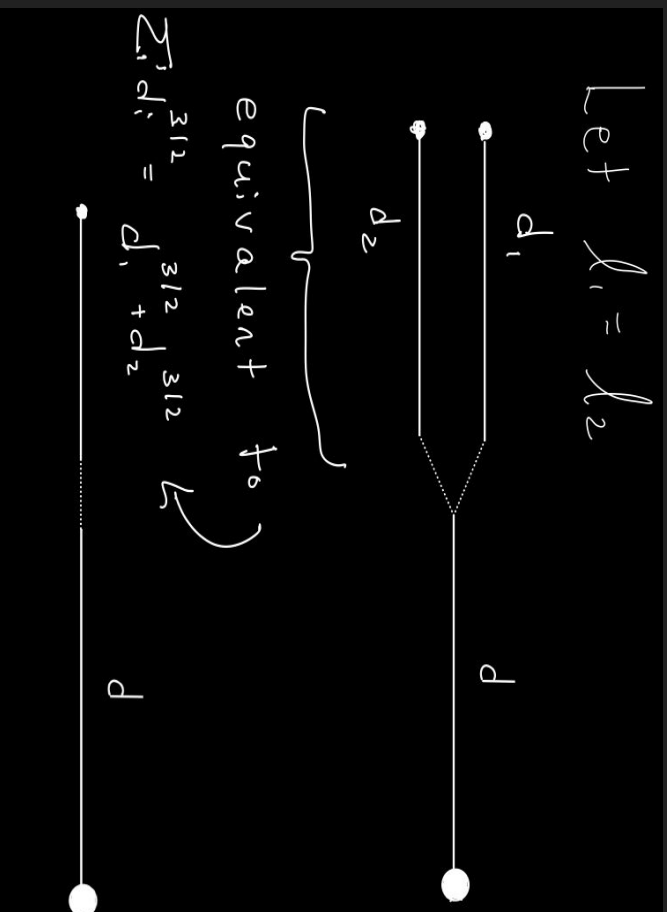


equivalent to

$$\sum_i d_i^{3/2} = d_1^{3/2} + d_2^{3/2}$$

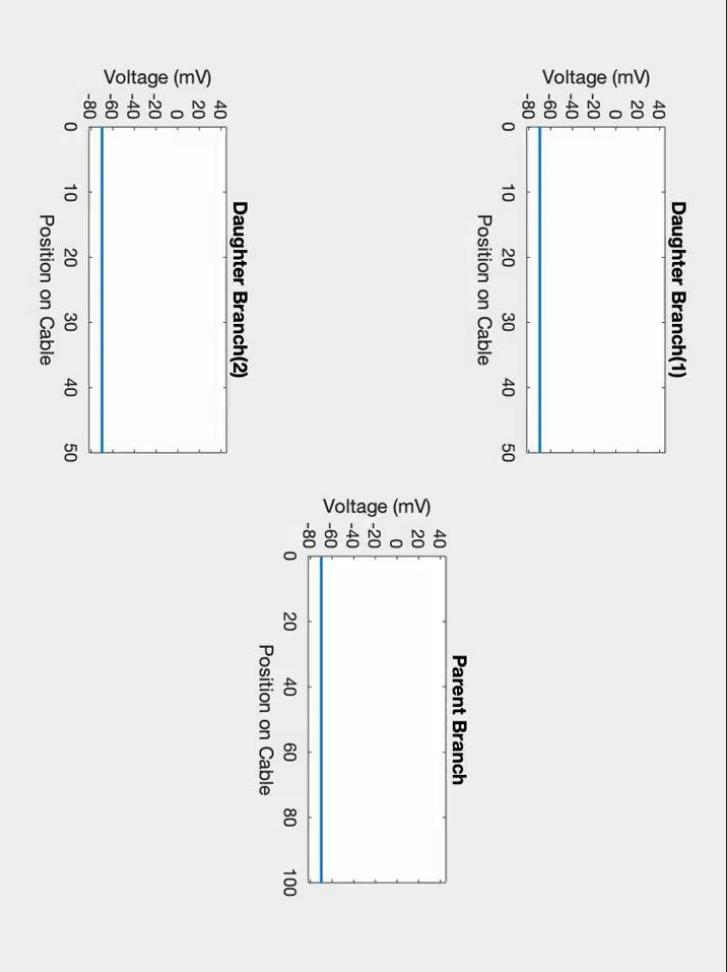
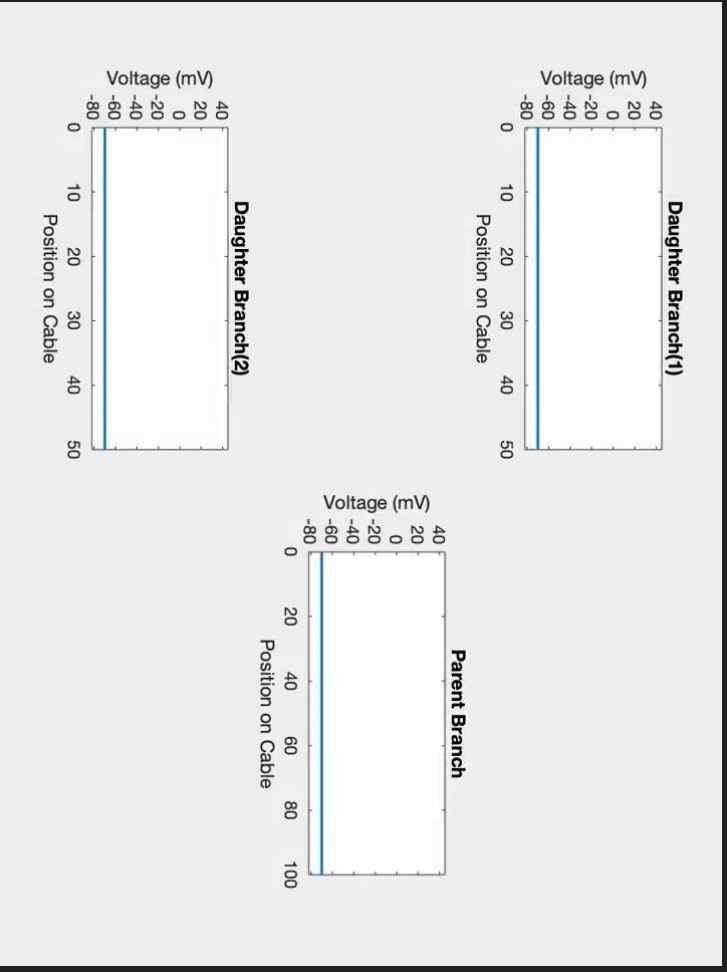
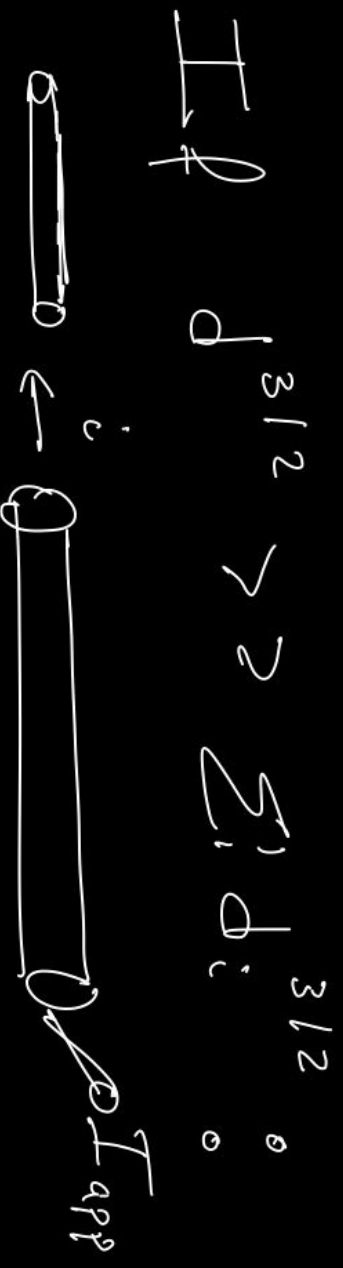


Rall's Equivalent Cylinder (1962)



- Since the electrotonic lengths of the daughters is the the same, we can reduce the branched case.
- This should ensure that propagation will follow “all or none” behavior.

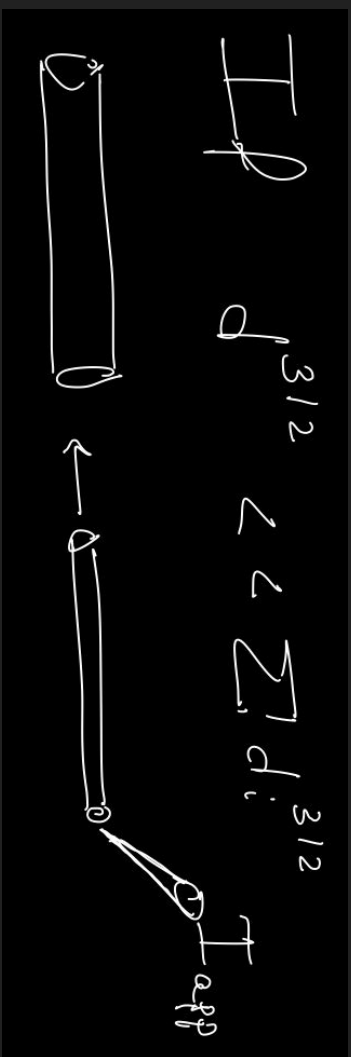
Case 1:



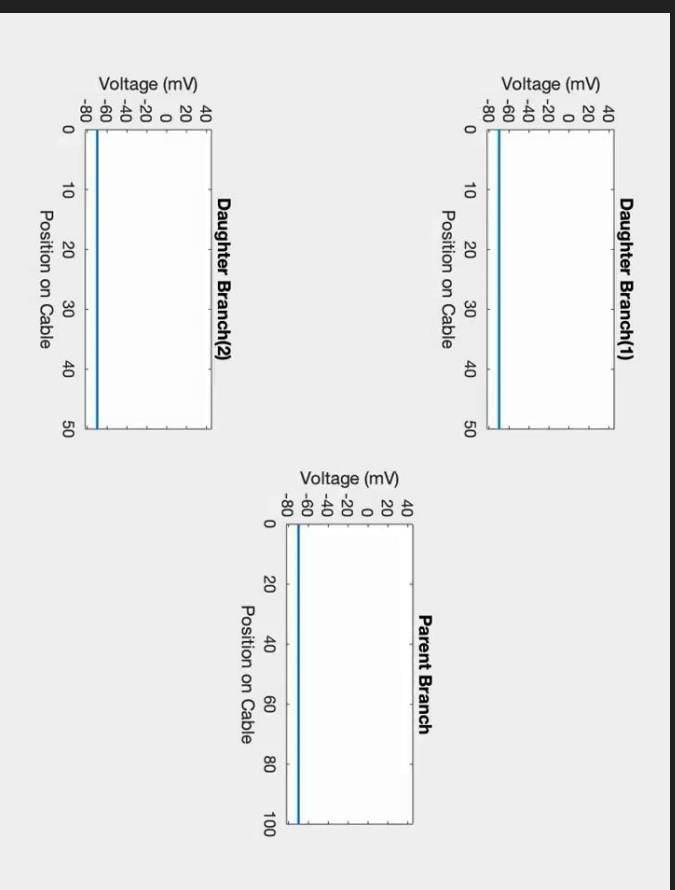
Check that solutions on each branch are the same :

```
>> v(401:600)-v(201:400)
ans =
Columns 1 through 23
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 24 through 46
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 47 through 69
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 70 through 92
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 93 through 115
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 116 through 138
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 139 through 161
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 162 through 184
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
Columns 185 through 200
0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0 0
```

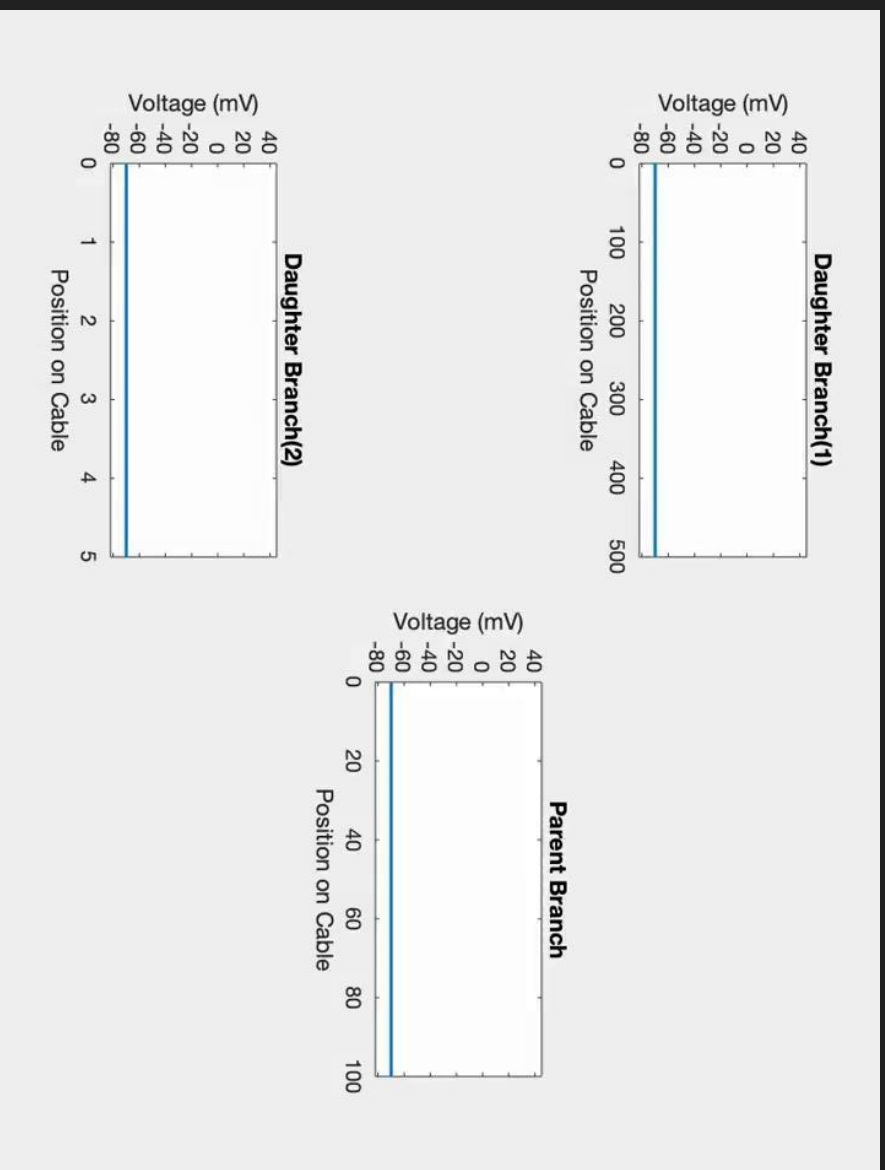
Case 2:



Too much membrane area to cover... propagation will FAIL



Differential Propagation: Unequal Electrotonic Lengths



References

- [1] Peskin, Charles S. “Mathematical Aspects of Neurophysiology .” Hodgkin Huxley Equations. 2000.
- [2] Rinzel J, Rall W. Transient response in a dendritic neuron model for current injected at one branch. *Biophys J.* 1974;14(10):759-790.
doi:10.1016/S0006-3495(74)85948-5

Link to the google slides for access to the videos:

<https://docs.google.com/presentation/d/1nVklidOWmzgTDwWfFoqgFYqK1C7cdr1lv82MetPq1-4/edit?usp=sharing>